



## Two-level system in an in-plane rotating magnetic field

Time dependent Hamiltonian  $H = \vec{\Omega}(t) \cdot \vec{\sigma}$   $\vec{\Omega} = \Omega(\cos \omega t, \sin \omega t, 0)^T$

Adiabatic limit, instantaneous eigenstate  $|\psi_{ad}(t)\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-iat} \\ 1 \end{pmatrix}$   
 → Berry phase [1]

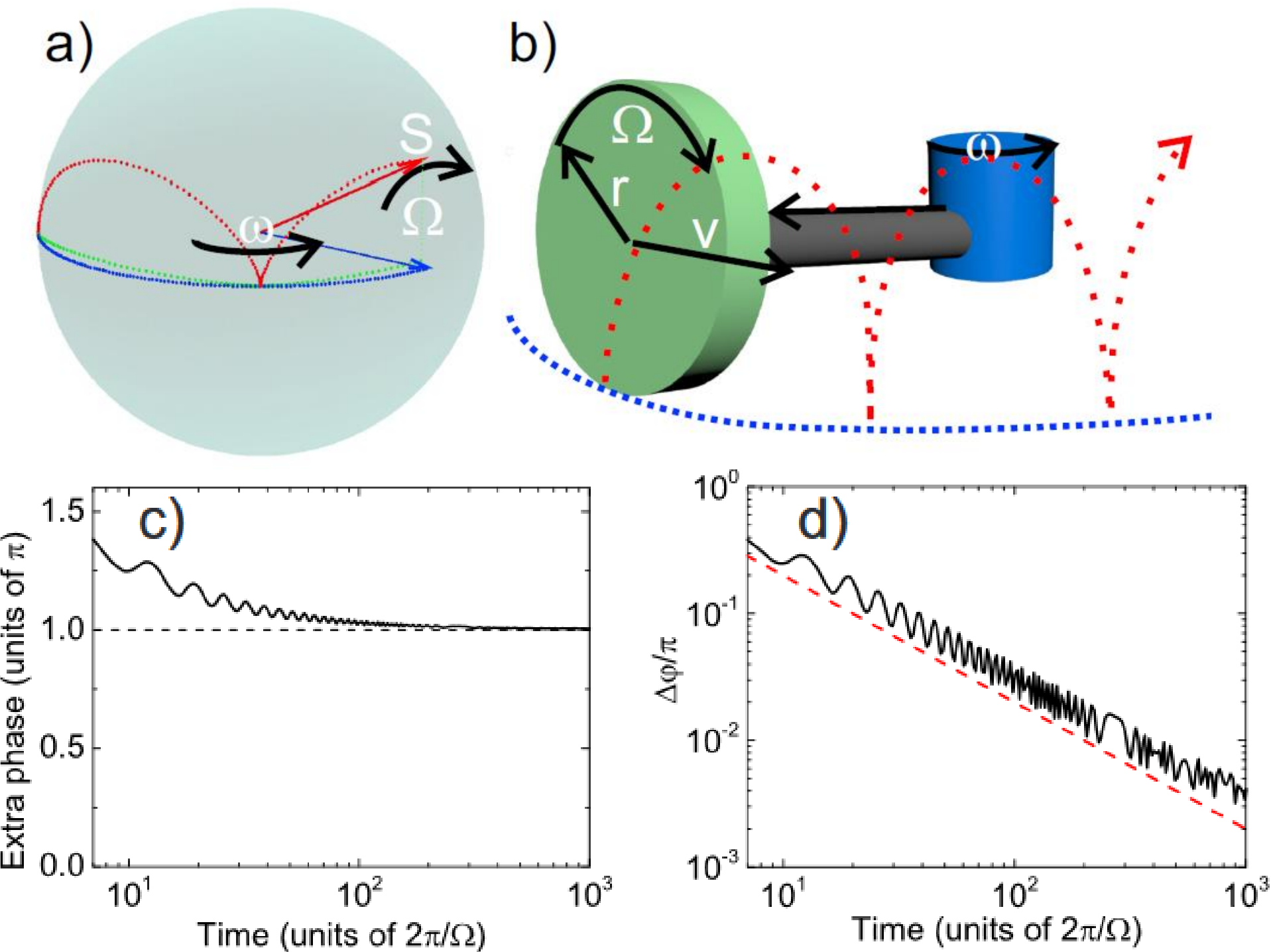
for 1 full rotation of  $\vec{\Omega}$ :  $\varphi_B = i \int_0^T \langle \psi_{ad} | \frac{\partial}{\partial t} | \psi_{ad} \rangle dt = \pm \pi$

Due to finite time experiment, perfect adiabaticity never reached (Cycloidal trajectory of the spin)

→ Correction to the geometric phase (see also [2])  
 → Non-adiabatic fraction

$$\varphi_{cor} = \pi \left( 1 + 2 \frac{\omega}{\Omega} \right)$$

$$\Delta \varphi_{cor} / \varphi_B = 4\pi / \Omega T$$



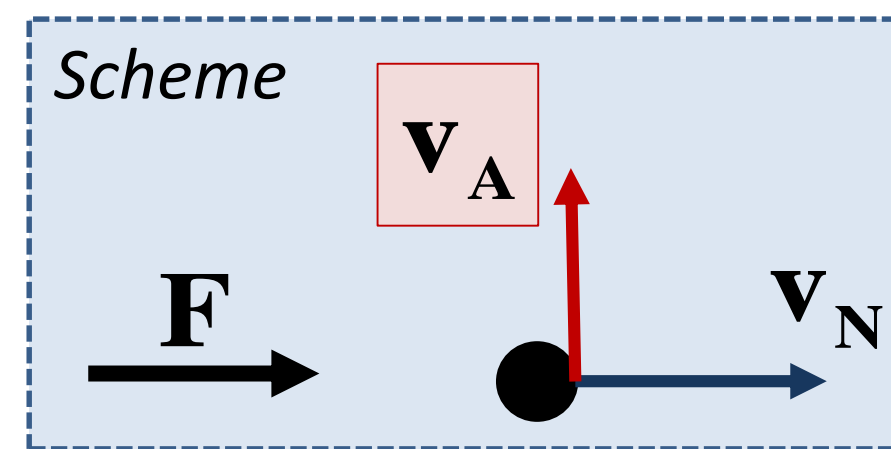
## Anomalous Hall effect [3,4]

Semiclassical equations of motion for wavepacket center of mass

$$\hbar \frac{\partial \mathbf{k}}{\partial t} = \mathbf{F}, \quad \text{External constant force}$$

$$\hbar \frac{\partial \mathbf{r}}{\partial t} = \frac{\partial E_n}{\partial \mathbf{k}} - \hbar \frac{\partial \mathbf{k}}{\partial t} \times \mathbf{B}_n(\mathbf{k}) \quad (1)$$

**Anomalous velocity**



Valid in the adiabatic limit

## Nonadiabatic correction to the semiclassical equations for a 2-band system

$$\Omega_{NA,ij}(\mathbf{k}) = \left( 1 + \sum_{l,m} \frac{g_{lm}}{(\epsilon_0 - \epsilon_l)^2} \frac{\partial k_l}{\partial t} \frac{\partial k_m}{\partial t} \right) \Omega_{ij}(\mathbf{k}) \quad \hbar \frac{\partial \mathbf{r}}{\partial t} = \frac{\partial \tilde{\epsilon}}{\partial \mathbf{k}} - \hbar \frac{\partial \mathbf{k}}{\partial t} \times \mathbf{B}_{NA} \quad (2)$$

$\Omega_{ij}$ : adiabatic Berry curvature

$g_{ij}$ : quantum metric component

$$\mathbf{B}_{NA} = (\Omega_{NA,yz}, \Omega_{NA,zx}, \Omega_{NA,xy})^T$$

Time evolution of a wavepacket in an energy band using both real and imaginary part of QGT computed with the instantaneous Hamiltonian

Valid on geodesics in momentum space, where 1<sup>st</sup> order corrections are zero

## Quantum geometric tensor (QGT) in k space

Introduced in 1980 to define metric in arbitrary parameter space in QM [5]

$$T_{ij}^{(n)}(\mathbf{k}) = \langle \partial_{k_i} u_n | \partial_{k_j} u_n \rangle - \langle \partial_{k_i} u_n | u_n \rangle \langle u_n | \partial_{k_j} u_n \rangle = g_{ij}^{(n)} - i \frac{\Omega_{ij}^{(n)}}{2}$$

→ Contains both quantum metric + Berry curvature

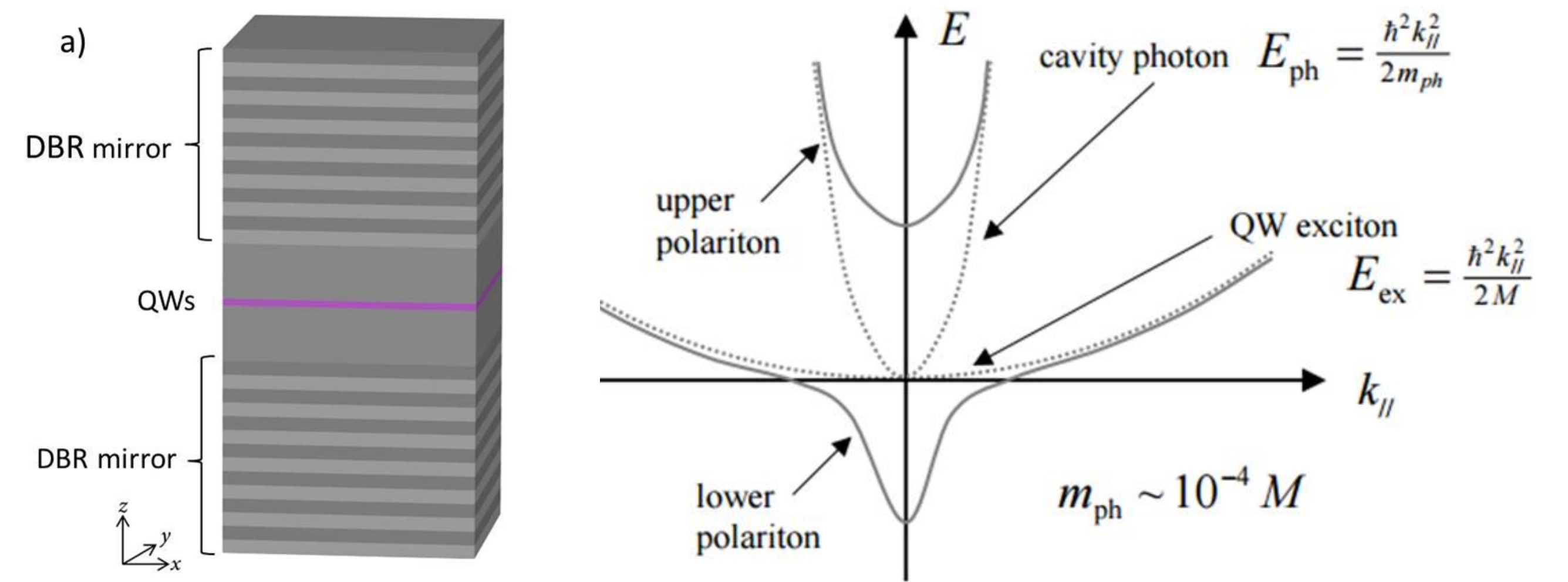
- Real part = metric  $ds^2 = \Re[T_{ij}] dk_i dk_j = g_{ij}^{(n)} dk_i dk_j$

$$ds^2 = 1 - |\langle u_n(\mathbf{k}) | u_n(\mathbf{k} + \delta \mathbf{k}) \rangle|^2$$

- Imaginary part = Berry curvature

$$\Omega_{ij}^{(n)} = -2\Im[T_{ij}^{(n)}] \quad \mathbf{B}_n = (\Omega_{yz}^{(n)}, \Omega_{zx}^{(n)}, \Omega_{xy}^{(n)})^T$$

## Cavity exciton-polaritons: strong light-matter coupling

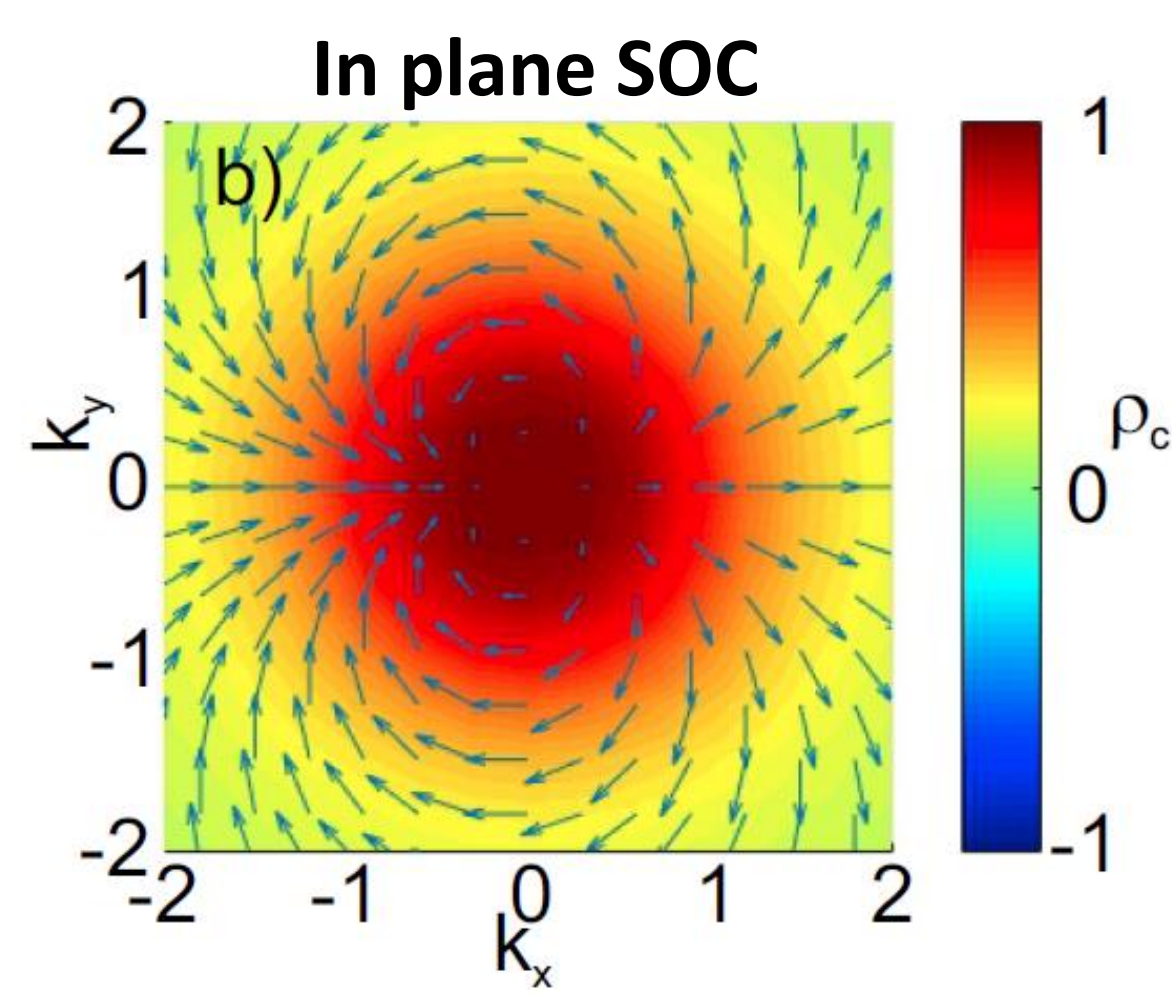
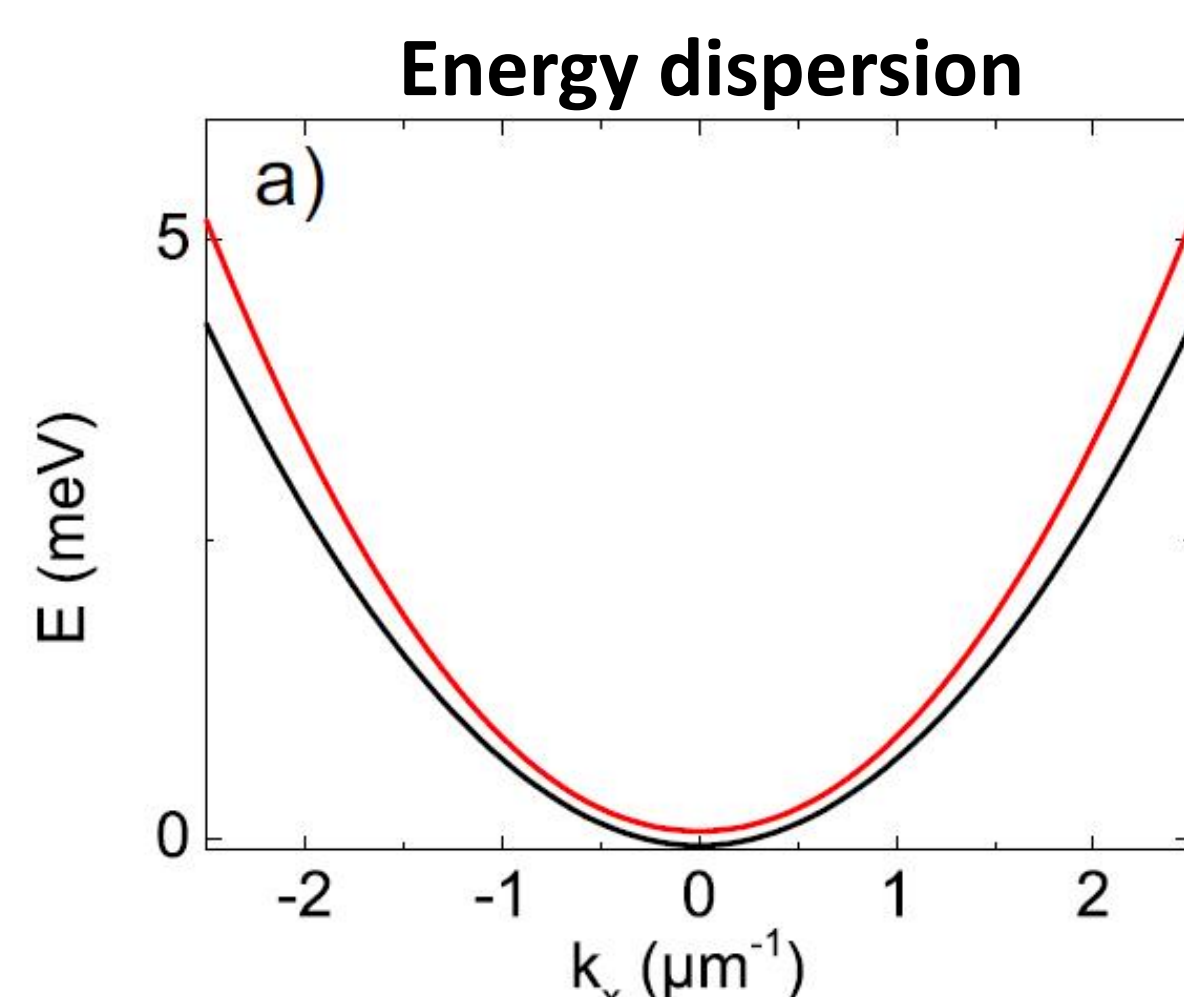


→ 2D neutral bosonic quasi-particles

## Quantum geometry of planar cavity polaritons

- LPB parabolic approximation
- TE-TM SOC + Zeeman field
- Effective 2-band system

$$H = \begin{pmatrix} \frac{\hbar^2 k^2}{2m^*} + \Delta & \beta k^2 e^{2i\phi} \\ \beta k^2 e^{-2i\phi} & \frac{\hbar^2 k^2}{2m^*} - \Delta \end{pmatrix}$$



### QGT components

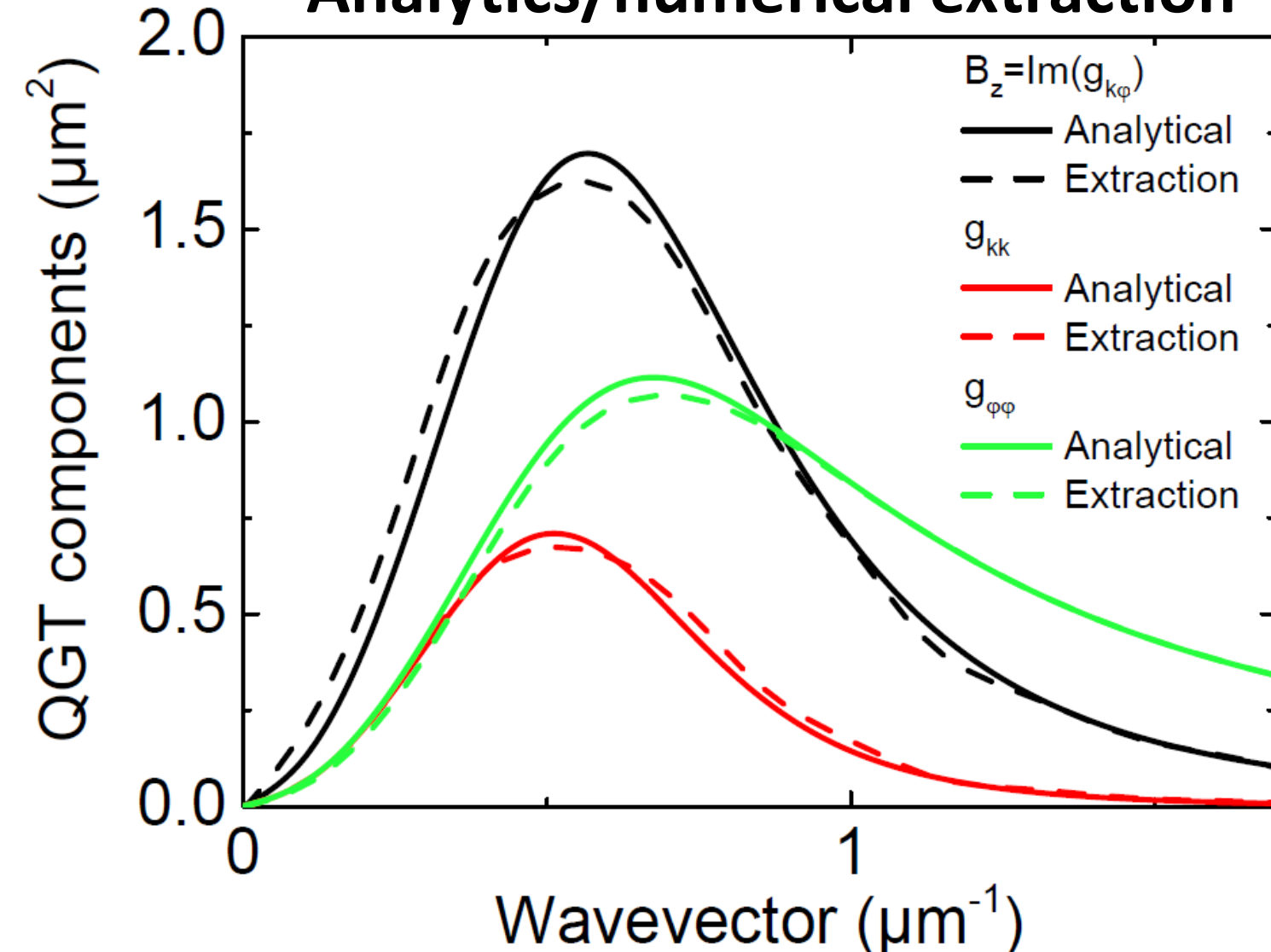
$$g_{kk} = \frac{\Delta^2 k^2 \beta^2}{(\Delta^2 + \beta^2 k^4)^2}$$

$$g_{\phi\phi} = \frac{k^2 \beta^2}{\Delta^2 + \beta^2 k^4}$$

$$g_{k\phi} = g_{\phi k} = 0$$

$$\mathbf{B} = \frac{2\Delta k^2 \beta^2}{(\Delta^2 + k^4 \beta^2)^{3/2}} \mathbf{e}_z$$

### Analytics/numerical extraction

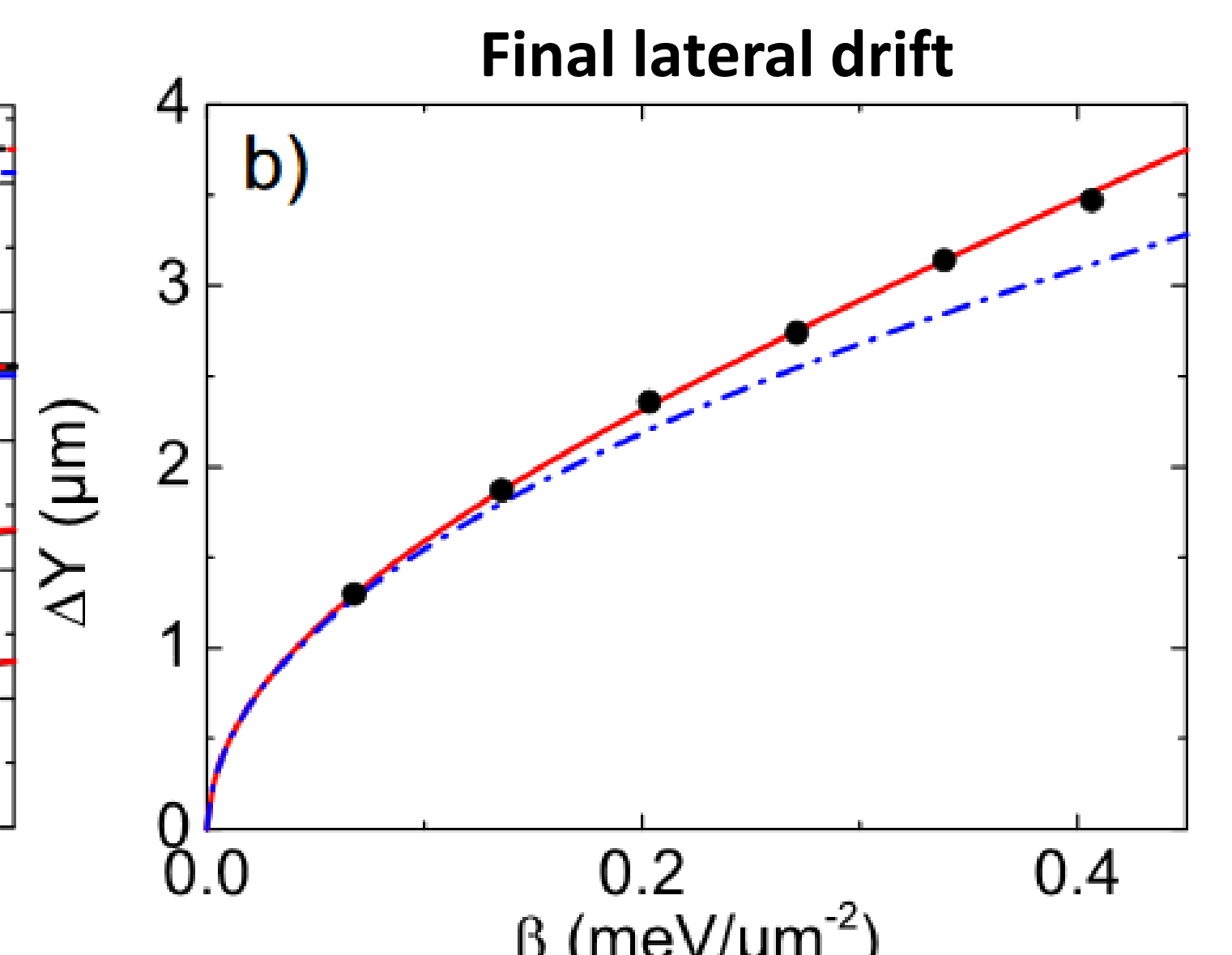
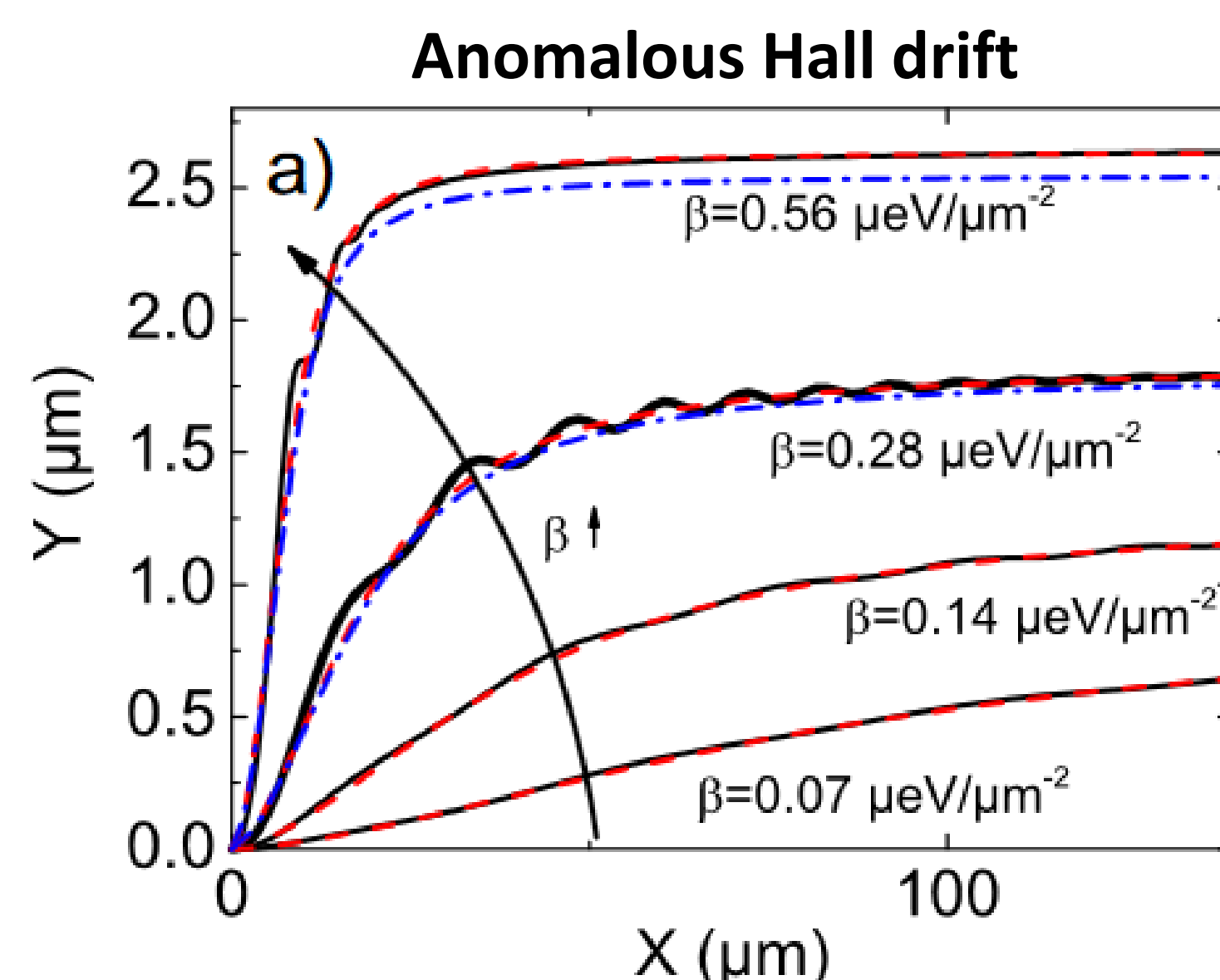


## Time evolution of a polariton wavepacket: Numerics/Analytics

Full numerical simulation using spinor Schrödinger eq.

Wedge potential  $U(x) = -Fx$

$$i\hbar \frac{\partial}{\partial t} \psi_{\pm} = -\frac{\hbar^2}{2m} \Delta \psi_{\pm} - \frac{i\hbar}{2\tau} \psi_{\pm} \pm \Delta \psi_{\pm} + \beta \left( \frac{\partial}{\partial x} \mp i \frac{\partial}{\partial y} \right) \psi_{\mp} + U(x) \psi_{\mp} + \hat{P}$$



Analytics: - without non-adiabatic correction (blue-dashed-eq.(1))

- with non-adiabatic correction (red-eq.(2))

Full Schrödinger numeric (black)

## Conclusions

- Berry curvature + quantum metric = QGT
- Directly measurable in planar cavities and cavity lattices (see [6])
- Both components important for realistic (nonadiabatic) AHE trajectory

## References

[1] M. V. Berry, *Proc. R. Soc. London, Ser. A* 392, 4557 (1984)  
 [2] K. Yu. Bliokh, *Phys. Lett. A* 372 204-9 (2008)  
 [3] R. Karplus, J. M. Luttinger, *Phys. Rev.* 95, 1154 (1954)

[4] G. Sundaram, Q. Niu, *Phys. Rev. B* 59, 14915 (1999)  
 [5] J. Provost, G. Vallee, *Comm. Math. Phys.* 76, 289 (1980)  
 [6] O. Bleu et al, *Phys. Rev. B* 97, 195422 (2018)