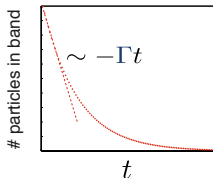
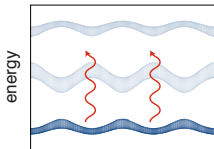
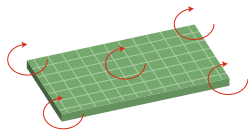


# Probing Quantum Geometry by “Heating”: From Quantized Circular Dichroism to Topological Order

Nathan Goldman

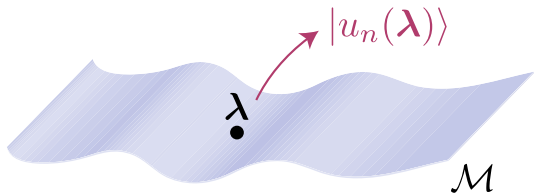


Aussois, 27th November 2018

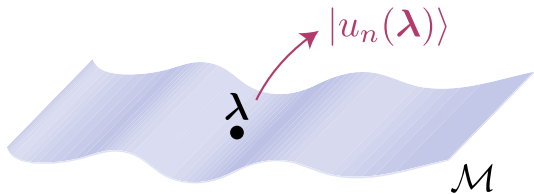
## Collaborations :

- Theory
  - N. R. Cooper (Cambridge)
  - A. Dauphin (ICFO)
  - A. G. Grushin (Grenoble)
  - T. Ozawa (Riken)
  - C. Repellin (MIT)
  - G. Palumbo and D. T. Tran (ULB)
  - P. Zoller (IQOQI)
- Experiments (group of K. Sengstock in Hamburg)
  - L. Asteria
  - N. Fläschner
  - B. S. Rem
  - K. Sengstock
  - M. Tarnowski
  - C. Weitenberg

- System described by Hamiltonian  $\hat{H}(\boldsymbol{\lambda})$  with  $\boldsymbol{\lambda} \in \mathcal{M}$
- Local eigenstates :  $\hat{H}(\boldsymbol{\lambda})|u_n(\boldsymbol{\lambda})\rangle = \varepsilon_n(\boldsymbol{\lambda})|u_n(\boldsymbol{\lambda})\rangle$



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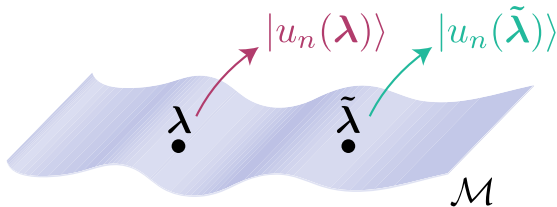


- Projection on  $n$ th band  $\varepsilon_n$  (adiabatic evolution) :

$$|\psi(t)\rangle = e^{i\theta(t)}|u_n[\boldsymbol{\lambda}(t)]\rangle, \quad \forall \text{ time } t$$

→ geometric structure related to the Hilbert space !

- **Geometry** : (a) distance between states



- The quantum (Fubini-Study) metric :

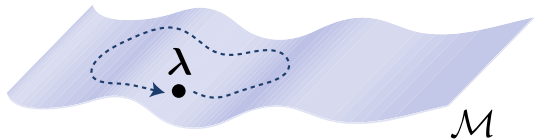
$$g_{\mu\nu} d\lambda^\mu d\lambda^\nu = 1 - |\langle u_n(\lambda) | u_n(\lambda + d\lambda) \rangle|^2$$

[Manifestations in condensed-matter, quantum information, . . . ]

[see review : Kolodrubetz, Polkovnikov et al. Phys. Rep. 697 (2017)]

- **Geometry** : (b) the curvature (geometric phase)

$$|u_n(\boldsymbol{\lambda})\rangle \longrightarrow |u_n(\boldsymbol{\lambda})\rangle \exp\left(i \int_{\Sigma} \boldsymbol{\Omega}^{(n)} \cdot d\mathbf{S}\right)$$



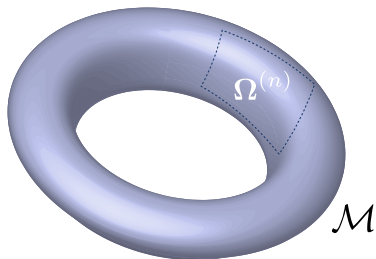
- The Berry curvature (in band  $\varepsilon_n$ ):

$$\boldsymbol{\Omega}^{(n)} = \nabla_{\boldsymbol{\lambda}} \times \left( i \langle u_n | \nabla_{\boldsymbol{\lambda}} u_n \rangle \right)$$

- “Aharonov-Bohm” phase in parameter space :

→ Berry curvature  $\equiv$  “magnetic field” in  $\boldsymbol{\lambda}$  space

- **Topology** : Gauss-Bonnet in quantum physics



- The Chern number of the  $n$ th band  $\varepsilon_n$  :

$$\nu^{(n)} = \frac{1}{2\pi} \int_{\mathcal{M}} \Omega^{(n)} \cdot d\mathbf{S} \in \mathbb{Z}$$

- Total “flux” of the field  $\Omega^{(n)}$  through  $\mathcal{M}$

→  $\nu^{(n)} \equiv$  “monopole” charge in  $\lambda$  space

- **Metric and curvature united** : the geometric tensor

$$G_{\mu\nu}^{(n)} = \sum_{m \neq n} \frac{\langle u_n | \partial_\mu \hat{H} | u_m \rangle \langle u_m | \partial_\nu \hat{H} | u_n \rangle}{(\varepsilon_n - \varepsilon_m)^2} \quad (\text{in band } \varepsilon_n)$$

- Quantum metric :  $g_{\mu\nu}^{(n)} = \mathcal{R} \left( G_{\mu\nu}^{(n)} \right)$
- Berry curvature :  $\Omega_{\mu\nu}^{(n)} = -2 \mathcal{I} \left( G_{\mu\nu}^{(n)} \right)$  [note :  $(\Omega)^\alpha = \epsilon^{\alpha\mu\nu} \Omega_{\mu\nu}$ ]

**Universal method to extract  $G_{\mu\nu}$  ?**

[beyond state tomography ; see Fläschner et al. Science '16]



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- Prepare system at  $\lambda = \lambda^0$  in eigenstate  $|u_0(\lambda^0)\rangle$

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- Measure the excited fraction over time

$$n_{\text{ex}}(t) = \frac{2\pi t}{\hbar} \sum_{n \neq 0} \left| \frac{\mathcal{E}}{\hbar\omega} \langle n | \partial_\mu \hat{H}(\lambda^0) | 0 \rangle \right|^2 \delta(\varepsilon_n - \varepsilon_0 - \hbar\omega)$$

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- Integrate the excited rate  $\Gamma = n_{\text{ex}}(t)/t$  over frequency

$$\Gamma^{\text{int}} = \int_0^\infty \Gamma(\omega) d\omega = \frac{2\pi\mathcal{E}^2}{\hbar^2} g_{\mu\mu}(\lambda^0)$$

## A general protocol (2)

- If one modulates two parameters (say  $\lambda_\mu$  and  $\lambda_\nu$ )

$$\lambda_\mu(t) = \lambda_\mu^0 + 2(\mathcal{E}/\hbar\omega) \cos(\omega t)$$

$$\lambda_\nu^\pm(t) = \lambda_\nu^0 \pm 2(\mathcal{E}/\hbar\omega) \cos(\omega t)$$

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- Measure the differential integrated rate

$$\Delta\Gamma^{\text{int}} = \int_0^\infty \Gamma^+(\omega) - \Gamma^-(\omega) d\omega = \frac{8\pi\mathcal{E}^2}{\hbar^2} g_{\mu\nu}(\boldsymbol{\lambda}^0)$$

→ one reconstructs the full quantum metric !

## A general protocol (3)

- If one further changes the relative phase

$$\lambda_{\mu}(t) = \lambda_{\mu}^0 + 2(\mathcal{E}/\hbar\omega) \cos(\omega t)$$

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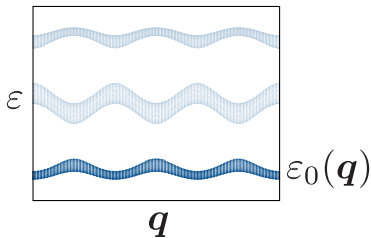
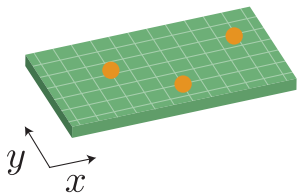
- Measure the differential integrated rate

$$\Delta\Gamma^{\text{int}} = \int_0^{\infty} \Gamma^{+}(\omega) - \Gamma^{-}(\omega) d\omega = \frac{4\pi\mathcal{E}^2}{\hbar^2} \Omega_{\mu\nu}(\boldsymbol{\lambda}^0)$$

→ one reconstructs the full Berry curvature !



## Application : Particles moving on a 2D lattice



- Parameters ( $\lambda$ ) are the quasi-momenta  $\mathbf{q} = (q_x, q_y)$
- Accessing  $\Omega_{xy}(\mathbf{q}^0)$  requires modulating

$$q_x(t) = q_x^0 + A \cos(\omega t)$$

$$q_y^\pm(t) = q_y^0 \pm A \sin(\omega t)$$

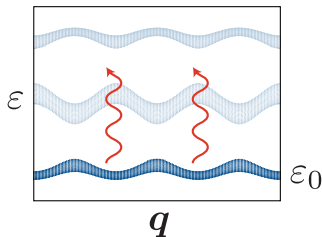
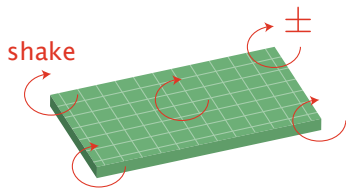
In practice ?

## Application : Particles moving on a 2D lattice

- Modulating the quasi-momenta ?

$$q_x(t) = q_x^0 + A \cos(\omega t)$$
$$q_y^\pm(t) = q_y^0 \pm A \sin(\omega t)$$

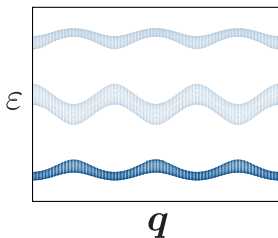
- Practice : **shake** the lattice circularly ( $\equiv$  circ. pol. light)



Geometry revealed through repopulation of the bands upon shaking !

## Application : Chern insulator

- **Filled** Bloch band with non-zero Chern number

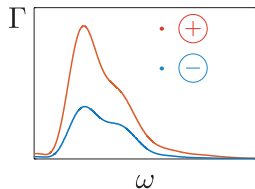
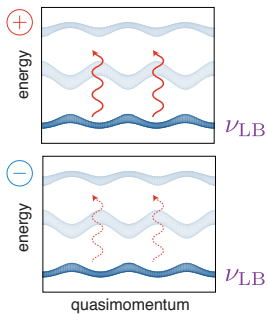
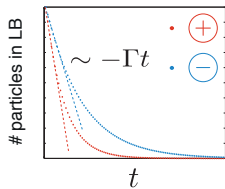
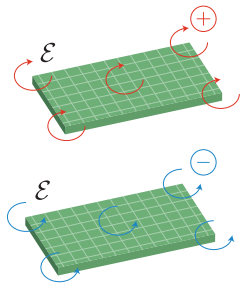


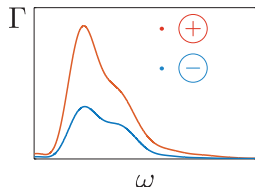
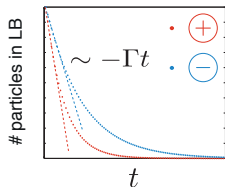
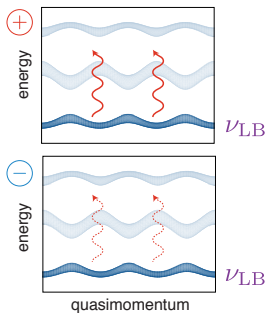
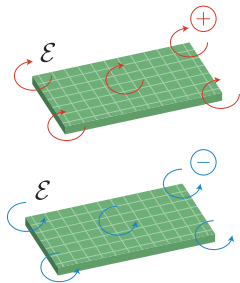
$$\nu = \frac{1}{2\pi} \int_{\text{FBZ}} \Omega \neq 0$$

- The protocol probes the **averaged** Berry curvature

$$\Delta\Gamma^{\text{int}} \sim \frac{1}{2\pi} \int_{\text{FBZ}} \Omega_{xy} dq_x dq_y = \nu \in \mathbb{Z}$$

→ a quantized dissipative response ?





## Quantization law for the rates

$$\Delta\Gamma^{\text{int}}/A_{\text{sys}} \equiv \frac{1}{2A_{\text{sys}}} \int_0^\infty (\Gamma^+ - \Gamma^-) d\omega = (\mathcal{E}/\hbar)^2 \times \nu_{\text{LB}}$$

## Quantized circular dichroism and relation to conductivity

- Circular dichroism and dissipative optical conductivity

$$\hbar\omega (\Gamma^+(\omega) - \Gamma^-(\omega)) = \mathcal{E}^2 (4A_{\text{syst}}) \times \sigma_1^{xy}(\omega)$$

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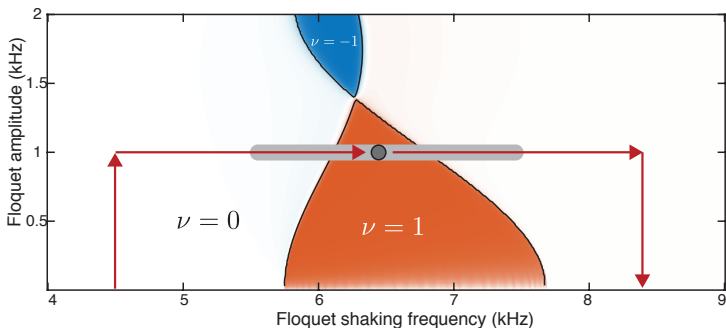
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- **Chern insulator case : quantized circular dichroism**

$$\sigma_{xy} = (1/h)\nu_{\text{LB}} \longrightarrow \Delta\Gamma^{\text{int}}/A_{\text{sys}} = (E^2/\hbar^2) \nu_{\text{LB}}$$

## Experimental demonstration of quantized circular dichroism

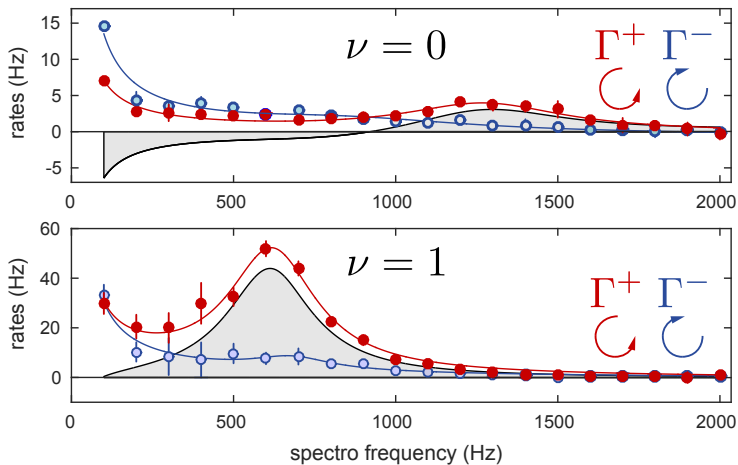
- $^{40}\text{K}$  atoms (fermions) in a honeycomb optical lattice
- Chern insulator is realized by Floquet engineering



- Excitation rates  $\Gamma^\pm$  are measured upon applying an additional circular shaking (orientation  $\pm$ )

# Experimental demonstration of quantized circular dichroism

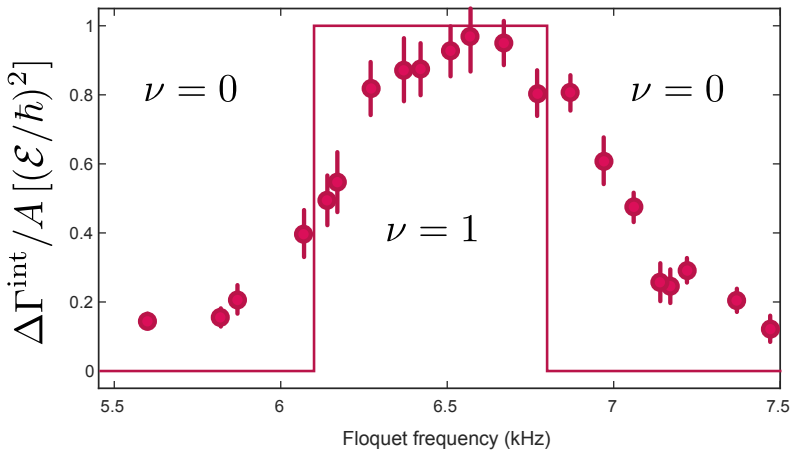
- Compare the rates in trivial and non-trivial regimes



# Experimental demonstration of quantized circular dichroism

- Quantized value reached deep inside topolog. regime

$$\Delta\Gamma^{\text{int}}/A_{\text{sys}} = (E^2/\hbar^2) \nu_{\text{LB}}^{\text{exp}}, \quad \text{with } \nu_{\text{LB}}^{\text{exp}} = 0.9(1)$$



## Other experimental results

- The experiment probes the **dissipative optical conductivity**

$$\Gamma^+(\omega) - \Gamma^-(\omega) = \mathcal{E}^2 \left( \frac{4A_{\text{sys}}}{\hbar\omega} \right) \times \sigma_1^{xy}(\omega)$$

→ first observation in topological quantum gases !

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- By shaking linearly, one probes the **quantum metric**

$$(\hbar/\mathcal{E})^2 (1/2\pi) \{ \Gamma_x^{\text{int}} + \Gamma_y^{\text{int}} \} = \overline{\text{Tr}[g_{\mu\nu}(\mathbf{k})]} \equiv \Omega_I,$$

→ exp. access to the **Wannier spread functional**

## Application : Fractional Chern insulators

- Universal relation for the differential integrated rate

$$\Delta\Gamma^{\text{int}}/A_{\text{syst}} = (2\pi\mathcal{E}^2/\hbar) \times \sigma_{xy}$$

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- Fractional Chern insulator :  $\sigma_{xy} = (1/h) \nu^{\text{MB}}$  ,  $\nu^{\text{MB}} \in \mathbb{Q}$

$$\Delta\Gamma^{\text{int}}/A_{\text{syst}} = (E^2/\hbar^2) \nu^{\text{MB}}$$



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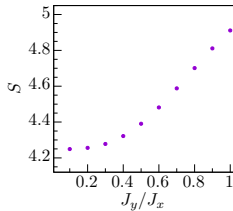
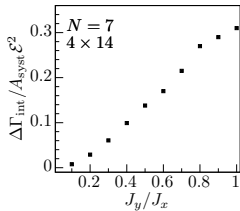
- Hardcore bosons on square lattice with flux  $n_{\Phi} = 1/4$

$N_{\text{part}}$	$N_{\text{supercell}}$	$\nu^{\text{MB}}$	$\Delta\Gamma^{\text{int}}/(A_{\text{syst}}\mathcal{E}^2)$
4	8	0.293	0.289
5	10	0.311	0.307
6	12	0.453	0.448
7	14	0.4638	0.45

→ converges towards  $\nu^{\text{MB}} = 0.5$  (**Laughlin state**)

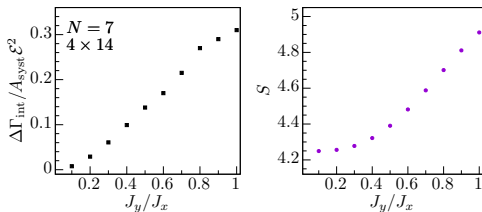
## Other results

- The dichroic signal distinguishes Laughlin from CDW

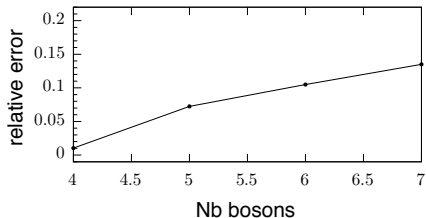


## Other results

- The dichroic signal distinguishes Laughlin from CDW



- Contributions from intra-band transitions are very small



→ good news for **band-mapping** !

# Conclusions

- Quantum geometry can be probed through excitation-rate measurements
- Quantized circular dichroism was revealed in ultracold atoms
- Detection schemes based on excitation-rate measurements are universal (wide range of applications)
- Circular dichroism can unambiguously detect topological order, e.g. fractional Chern insulators

## References

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Science Advances 3, e1701207 (2017)
- *Quantized Rabi Oscillations and Circular Dichroism in Quantum Hall Systems*,  
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- *Detecting fractional Chern insulators through circular dichroism*,  
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### Experimental results :

- *Measuring quantized circular dichroism in ultracold topological matter*,  
L. Asteria, D. T. Tran, T. Ozawa, M. Tarnowski, B. S. Rem, N. Fläschner,  
K. Sengstock, N. Goldman and C. Weitenberg,  
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### Not discussed but related :

- *Revealing tensor monopoles through quantum-metric measurements*,  
G. Palumbo and N. Goldman,  
Phys. Rev. Lett. 121, 170401 (2018) + [arXiv:1811.02434](#)